# A NEW EQUATION FOR THE HEIGHT EQUIVALENT TO A THEORETICAL PLATE

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#### INTRODUCTION

A fundamental system of differential equations for the mass transfer in a continuous column, under the assumption of a uniform fluid velocity, has been proposed and solved by LAPIDUS AND AMUNDSON<sup>1</sup>. VAN DEEMTER et al.<sup>2</sup> convert the solution into the familiar form of Gaussian distribution. The concept of HETP introduced from the plate theory<sup>3</sup>, is also dealt with by VAN DEEMTER et al.<sup>2</sup> to explain the band broadening and the column efficiency. They derive the following famous equation:

$$H = A + \frac{B}{u} + Cu \tag{1}$$

Here H is HETP, u the linear velocity of carrier gas and A, B and C are constants.

It is well known that the linear velocity of carrier gas varies with pressure gradient along the column and that the diffusion coefficient of gas is inversely proportional to the pressure. KAMBARA AND OHZEKI4 proposed a mathematical theory for the column, in which the kinetic role of pressure drop<sup>5</sup> and diffusion are considered, and derived the solution in a form of Gaussian distribution.

### THEORETICAL CONSIDERATIONS

The expression for the outlet concentration of solute in gas chromatography derived by KAMBARA AND OHZEKI<sup>4</sup> is shown by:

$$C(l,t) = \frac{p_o C_0 t_0}{p_i \sqrt{2\pi\sigma^2}} \exp\left\{-\frac{(t-t_R)^2}{2\sigma^2}\right\}$$
(2)

with

$$\sigma^2 = \sigma_D^2 + \sigma_\alpha^2 \tag{3}$$

$$\sigma_{D}^{2} = \frac{D^{*l}}{\beta^{2} u_{o}^{3} p_{o}} \left( \frac{p_{l}^{2} + p_{o}^{2}}{p_{o}^{2}} \right)$$
(4)

$$\sigma_{\alpha}^{2} = \frac{2F_{s}^{2}T_{l}}{\alpha F K^{2}} \tag{5}$$

J. Chromatog., 27 (1967) 33-39

$$t_{R} = \frac{T_{l}}{\beta}$$

$$T_{l} = \frac{4l^{2}(p_{l}^{3} - p_{o}^{3})}{3w(p_{l}^{2} - p_{o}^{2})^{2}}$$
(6)
(7)

and

$$v = qupF = \frac{wqF(p_{i}^{2} - p_{o}^{2})}{2l}$$
(8)

The significance of the symbols is summarized in Table I.

## TABLE I

#### LIST OF SYMBOLS

Symbol	Significance	Unit or dimension
с	concentration in gas phase	mole cm <sup>-3</sup>
Cs	same in stationary phase	mole cm <sup>-3</sup>
Co	initial concentration injected	mole cm <sup>-3</sup>
<b>D</b> *	diffusion constant under the unit atmosphere	cm² sec−1 atm
F	fraction by volume occupied by mobile phase	dimensionless
$F_S$	fraction by volume occupied by stationary phase	dimensionless
Ĥ	height equivalent to a theoretical plate, HETP	cm
K	distribution coefficient (= $C/C_s$ at equilibrium)	dimensionless
l	length of packing	cm
12	number of theoretical plates	dimensionless
<i>p</i> i	pressure of carrier gas at column inlet	atm $[ML^{-1} T^{-2}]$
Po	pressure of carrier gas at column outlet	atm [ML <sup>-1</sup> T <sup>-2</sup> ]
$\overline{q}$	proportionality factor in eqn. (8)	atm <sup>-1</sup> mole cm <sup>-3</sup>
to	duration of concentration pulse injected	sec
t <sub>R</sub>	<ul> <li>retention time of sample component</li> </ul>	SCC
$T_{l}$	passing time of carrier gas through the whole column length $l$	sec
26	linear gas velocity	cm sec <sup>-1</sup>
210	linear gas velocity at column outlet	cm sec <sup>-1</sup>
υ	gas flow rate per unit cross section of column	mole $cm^{-2} sec^{-1}$
w	permeability coefficient	$cm^2 sec^{-1} atm^{-1}$
æ	rate constant of dissolution or adsorption	sec <sup>-1</sup>
β	$I + (F_S/KF)$	dimensionless
$\sigma^{2}$	variance of peak	sec <sup>2</sup>
$\sigma_D^2$	contribution of diffusion term to variance	sec <sup>2</sup>
$\sigma_{\alpha}{}^2$	contribution of adsorption rate term to variance	Sec <sup>2</sup>

The HETP is given by:

$$H=\frac{l}{n}=l\cdot\frac{\sigma^2}{t_R^2}$$

÷

Hence, it can easily be derived that

:

$$H = \frac{9D^{*l}(p_i^4 - p_o^4)}{2w(p_i^3 - p_o^2)^3} + \frac{3wF_s^2\beta^2(p_i^2 - p_o^2)^2}{2\alpha F K^2 l(p_i^3 - p_o^3)}$$

J. Chromatog., 27 (1967) 33-39

(10)

(9)

34

The HETP is also related to  $u_0$ , the linear velocity of carrier gas at the column outlet, according to:

$$H = \frac{18D^{*l^{2}}}{w^{3}} \times \frac{u_{o}p_{o}(lu_{o}p_{o} + wp_{o}^{2})}{\left\{ \left( \frac{2lu_{o}p_{o}}{w} + p_{o}^{2} \right)^{3/2} - p_{o}^{3} \right\}^{2}} + \frac{6F_{s}^{2}\beta^{2}l}{\alpha FK^{2}w} \times \frac{(u_{o}p_{o})^{2}}{\left\{ \left( \frac{2lu_{o}p_{o}}{w} + p_{o}^{2} \right)^{3/2} - p_{o}^{3} \right\}}$$
(11)

If the condition:

$$p_i^2 \gg p_o^2 \tag{12}$$

holds, one has:

$$H = \frac{9D^*}{4u_o p_o} + \frac{3wF_s^2\beta^2}{2\alpha FK^2l} \sqrt{\frac{2u_o p_o l}{w}}$$
(13)

Thus, in the case where the carrier gas flow rate is very high, the simplified form for HETP is shown by:

$$H = \frac{B}{u_o} + C\sqrt{u_o} \tag{14}$$

or

$$H = \frac{B'}{v} + C' \sqrt{v} \tag{15}$$

EXPERIMENTAL

The gas chromatograph used was a Hitachi product, type KGL-2A, equipped with a thermal conductivity detector. The column, 2 m in length and 0.4 cm in diameter, was a copper tube packed with a stationary phase of 25 % DNP on 60-80 mesh Shimalite W. The difference between the inlet and outlet pressures was measured by means of a pressure gauge and it was confirmed that under the conditions of operation the outlet pressure was always kept constant at atmospheric within *ca*.  $\pm$  1%, as shown in Table II. Helium carrier gas flow rate was measured with a soap film meter and corrected for the saturated water vapour pressure. The column temperature was maintained at 77  $\pm$  0.1°, at which temperature the chromatogram of benzene can be considered to obey approximately the normal distribution law. The number of plates and the HETP for benzene were found to be independent of the sample size with amounts smaller than 0.8  $\mu$ l and so they were determined for 0.6  $\mu$ l samples.

#### **RESULTS AND DISCUSSION**

Values of the flow rate and the HETP obtained from chromatograms of benzene under varying inlet pressures are summarized in Table II.

The flow rate v, expressed in mole sec<sup>-1</sup> cm<sup>-2</sup>, is plotted against  $(p_i^2 - p_0^2)$  in Fig. 1. The results shows a good agreement with eqn. (8).

### TABLE II

VARIATIONS OF CARRIER GAS FLOW RATE AND HETP WITH PRESSURE DIFFERENCE

Pressure difference, Þi — Þo (atm)	Outlet pressure, po (atm)	Carrier gas flow rate (10 <sup>-5</sup> mole cm <sup>-2</sup> sec <sup>-1</sup> )	HETP (mm)
0.39	0,990	0.65	1.32
0.59	0.992	1.09	0.98
0.78	0.990	1.54	0,88
0.97	0.988	2.07	0.85
1.17	0.988	2.64	0.87
1.36	0.983	3.24	o.88
1.55	0.983	3.96	0.94
1.75	0.980	4.68	0.99
1.94	0.980	5.41	1.06
2.14	0.979	6.34	1.15
2.33	0.982	7.18	1,21
2.52	0.986	8.17	1.30
2.72	0.979	9.12	1.38
2.91	0.975	10.21	1.44



Fig. 1. The absolute flow rate of the helium carrier gas plotted against the difference between the squares of inlet and outlet pressures as shown by eqn. (8).

The plot of HETP against v is shown in Fig. 2. As the flow rate increases, the plot becomes slightly convex as it curves upwards and deviates from the linearity predicted by the original VAN DEEMTER theory. Such a tendency has also been observed recently by LITTLEWOOD<sup>6</sup>.

J. Chromatog., 27 (1967) 33-39

The plot shown in Fig. 3 forcibly demonstrates the theoretical prediction derived from eqn. (10), which holds generally for a wide interval of pressure difference.

Next, the plot of log H against log v, shown in Fig. 4, indicates most clearly and definitely the essential validity of the present theory at higher flow rates. The slope of this straight line, which is evidently 0.5, reveals that HETP is proportional to  $\sqrt{v}$  or  $\sqrt{u_0}$  in this range and not to v or  $u_0$  itself.



Fig. 2. The HETP plotted against the absolute flow rate of the carrier gas. Sample: benzene. Column: 2 m long, 0.4 cm diam., 25 % DNP on 60-80 mesh Shimalite W at 77°. The vertical lines denote the 95 % confidence interval.



Fig. 3. Verification of the inlet and outlet pressure dependencies of the HETP as shown by eqn. (10).

In Fig. 5, the product Hv is plotted against  $v^{3/2}$ . The linear relationship at higher flow rates also proves the correctness of the present new HETP equation.

The constant term A in the original VAN DEEMTER equation (I) has been introduced under the assumption that the effective diffusion coefficient is an algebraic sum of the true diffusion coefficient multiplied by the labyrinth factor and the eddy diffusivity. If one regards that the former is independent of the linear velocity, while the latter is proportional to it, the term A appears in the HETP equation. In some earlier investigations, *e.g.* by BOHEMEN AND PURNELL<sup>7</sup>, it was, however, reported that the term A showed negative values indicating the weakness of the above assumption.



Fig. 4. Double-logarithmic plot of HETP against the absolute flow rate of carrier gas. The slope of the linear part in the higher flow rate region clearly shows the correctness of the revised HETP expression given by eqns. (14) and (15).



Fig. 5. Plot of the product of HETP and the absolute flow rate v of carrier gas against  $v^{3/2}$  as predicted by the revised eqn. (15) for HETP.

SUMMARY

On the basis of our previous kinetic theory, in which the pressure dependencies of the diffusion coefficient in the gas phase and the linear velocity distribution along

J. Chromalog., 27 (1967) 33-39

the column are taken into account, a new expression for the gas chromatographic HETP is derived theoretically. Experimental results definitely confirm the essential validity of the present revised equation for HETP. It is concluded that in the VAN DEEMTER equation the constant term A should disappear and the term  $Cu_0$  should be proportional not to the linear velocity  $u_0$  but to the square root of it.

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J. Chromalog., 27 (1967) 33-39